

**G. Eigen (Bergen)** 

2nd physics of the B factories workshop KEK, May 18, 2010 On behalf of M. Bona, G. Eigen, R. Itoh and E. Kou

# Chapter Outline

- Section: Introduction and goals 2p
- Section: Methodology
  - Subsection: CKMfitter 2p
  - Subsection: UTfit 2p
  - Subsection: Scanning method 2p
- Section: Experimental Inputs
  - Subsection: B-factories results:  $\beta$ ,  $\alpha$  (which decays to consider),  $\gamma$ ,  $2\beta + \gamma$ ,  $V_{ub}$ ,  $V_{cb}$ ,  $\Delta m_d$ ,  $A^d_{SL}$ , B(B $\rightarrow \tau v$ ), radiative penguins (how to use them) 4p
  - Subsection: Non-B-factories results (briefly on their threatment):  $\varepsilon_{k}$ ,  $\Delta m_{s}$ ,  $A^{s}_{SI}$ , TD  $B_s \rightarrow J/\psi \phi$ ,  $\Delta \Gamma_s$  (with order calculation). 2-3p
    - Rather than having subsubsections we indicate in the table which are inputs for the SM fit and inputs for the BSM fits
- Section Theoretical Inputs
  - Subsection Derivation of hadronic observables 2p
  - Subsection Lattice QCD inputs 4p
- Benchmark models 5p
- Section Results from the global fits
- Section Global fits beyond the Standard Model 4p
  - Subsection New-physics parameterizations 4p
  - Subsection Operator analysis 2p
  - Section Conclusions 1-2p



total: 36-38 pages

# Motivation

The CKM matrix is specified by 4 independent parameters, in the Wolfenstein approximation they are  $\lambda$ , A,  $\overline{\rho}$ , and  $\overline{\eta}$ 

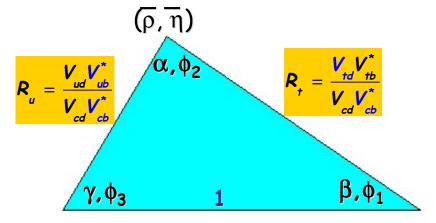
$$V_{CKM} = \begin{pmatrix} 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda + A^2\lambda^5(\frac{1}{2} - \rho - i\eta) & 1 - \frac{1}{2}\lambda^2 - \frac{1}{8}\lambda^4 - \frac{1}{2}A^2\lambda^4 & A\lambda^2 \\ A\lambda^3(1 - \overline{\rho} - i\overline{\eta}) & -A\lambda^2 + A\lambda^4(\frac{1}{2} - \rho - i\eta) & 1 - \frac{1}{2}A^2\lambda^4 \end{pmatrix} + O(\lambda^6)$$

Unitarity of the CKM matrix specifies relations among the parameters

e.g.  $V_{ud}V_{ub}^{\star} + V_{cd}V_{cb}^{\star} + V_{td}V_{tb}^{\star} = 0$ 

Combine measurements from the B and K systems to overconstrain the triangle

→ test if phase of CKM matrix is only source of CP violation





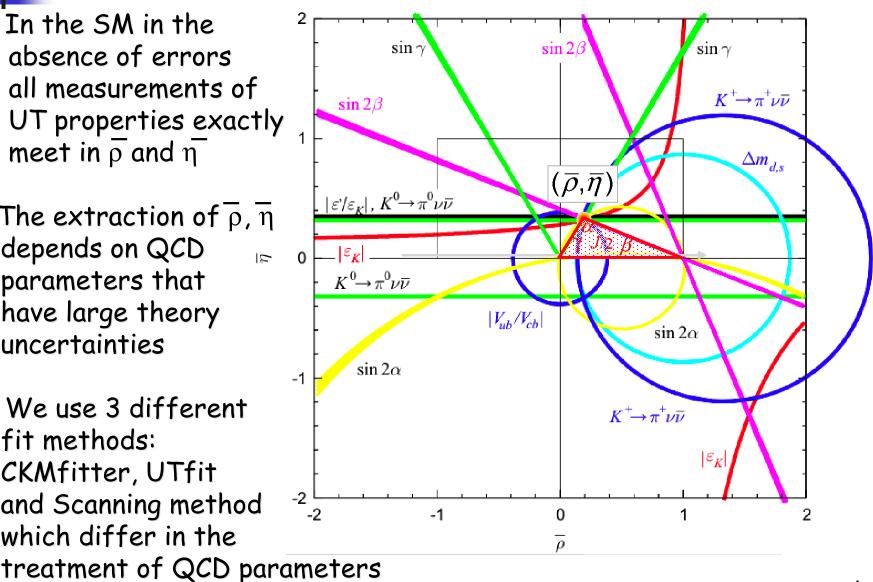
$$rg\left(-\frac{V_{td}V_{tb}^{\star}}{V_{ud}V_{ub}^{\star}}\right) \quad \beta, \phi_{1} = arg\left(-\frac{V_{cd}V_{cb}^{\star}}{V_{td}V_{tb}^{\star}}\right) \quad \gamma, \phi_{3} = arg\left(-\frac{V_{cd}V_{cb}^{\star}}{V_{td}V_{tb}^{\star}}\right)$$

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 $\alpha, \phi_2 = \alpha$ 

# **Motivation**

- In the SM in the absence of errors all measurements of UT properties exactly meet in  $\rho$  and  $\eta$
- The extraction of  $\overline{\rho}, \overline{\eta}$ depends on QCD 15 parameters that have large theory uncertainties
- We use 3 different fit methods: CKMfitter, UTfit and Scanning method which differ in the





# CKMfitter Methodology

- CKMfitter (Rfit) is a frequentist-based approach to the global fit of CKM matrix
- Likelihood function:

$$\mathcal{L}[\boldsymbol{y}_{mod}] = \mathcal{L}_{exp}[\boldsymbol{x}_{exp} - \boldsymbol{x}_{th}(\boldsymbol{y}_{mod})] \times \mathcal{L}_{th}[\boldsymbol{y}_{QCD}]$$

- $\rightarrow$  First term measures agreement between data,  $x_{exp}$ , and prediction,  $x_{th}$
- Second term expresses our present knowledge on QCD parameters
- $\rightarrow$  y<sub>mod</sub> are a set of fundamental and free parameters of theory (m<sub>t</sub>, etc)
- Minimize

and determine

$$\chi^{2}(\gamma_{mod}) \equiv -2\ln(\mathcal{L}[\gamma_{mod}])$$

$$\Delta \chi^{2}(\boldsymbol{y}_{mod}) = \chi^{2}(\boldsymbol{y}_{mod}) - \chi^{2}_{min;\boldsymbol{y}_{mod}}$$

where  $\chi^2_{\text{min};\text{ymod}}$  is the absolute minimum value of  $\chi^2$  function

- Separate uncertainties of QCD parameters into statistical ( $\sigma$ ) and non-statistical (theory) uncertainties ( $\delta$ )
- statistical uncertainties are treated like experimental errors with a Gaussian likelihood

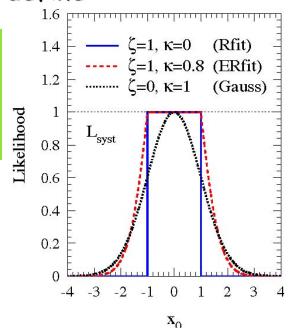
# Rfit Methodology

- Treatment of theory uncertainties ( $\delta$ ):
  - → If fitted parameter a lies within the predicted range  $x_0 \pm \delta x_0$  contribution to  $\chi^2$  is zero
  - → If fitted parameter a lies outside the predicted range x<sub>0</sub>±δx<sub>0</sub> the likelihood L<sub>th</sub>[ y<sub>QCD</sub>] drops rapidly to zero, define:

$$-2 \ln \mathcal{L}_{th}[\mathbf{x}_{0}, \kappa, \zeta] = \begin{cases} \mathbf{0}, & \forall \mathbf{x}_{0} \in \left[\overline{\mathbf{x}}_{0} \pm \zeta \delta \mathbf{x}_{0}\right] \\ \left(\frac{\mathbf{x}_{0} - \overline{\mathbf{x}}_{0}}{\kappa \delta \mathbf{x}_{0}}\right)^{2} - \left(\frac{\zeta}{\kappa}\right)^{2}, & \forall \mathbf{x}_{0} \notin \left[\overline{\mathbf{x}}_{0} \pm \zeta \delta \mathbf{x}_{0}\right] \end{cases}$$

- 3 different analysis goals
  - Within SM achieve best estimate of y<sub>th</sub>
  - Within SM set CL that quantifies agreement between data and theory
  - Within extended theory framework search for specific signs of new physics





# UTfit Methodology

UTfit is a Bayesian-based approach to the global fit of CKM matrix

- For M measurements  $c_j$  that depend on  $\rho$  and  $\eta$  plus other N parameters  $x_i$  the function  $f(\overline{\rho}, \overline{\eta}, x_1, \dots, x_N \mid c_1, \dots, c_M)$  needs to be evaluated by integrating over  $x_i$  and  $c_j$
- Using Bayes theorem one finds

$$\mathbf{f}(\overline{\rho},\overline{\eta},\mathbf{x}_{1},\ldots,\mathbf{x}_{N} \mid \mathbf{c}_{1},\ldots,\mathbf{c}_{M}) \propto \prod_{j=1,M} \mathbf{f}_{j}(\mathbf{c}_{j} \mid \overline{\rho},\overline{\eta},\mathbf{x}_{1},\ldots,\mathbf{x}_{N}) \prod_{i=1,N} \mathbf{f}_{i}(\mathbf{x}_{i}) \mathbf{f}_{0}(\overline{\rho},\overline{\eta})$$

where  $f_0(\overline{\rho}, \overline{\eta})$  is the a-priory probability for  $\overline{\rho}$  and  $\overline{\eta}$ 

• The output pdf for  $\rho$  and  $\eta$  is obtained by integrating over  $c_j$  and  $x_i$ 



# UTfit Methodology

Measurement inputs and all theory parameters are described by pdfs

- Errors are typically treated with a Gaussian model, only for  $B_k$ ,  $\xi$ and  $f_B \int B_B$  a flat distribution representing the theory uncertainty is convolved with a Gaussian representing the statistical uncertainty
- So if available, experimental inputs are represented by likelihoods
- The method does not make any distinction between measurement and theory parameters
- The allowed regions are well defined in terms of probability
   Allowed regions at 95% probability means that you expect the "true" value in this range with 95% probability
- By changing the integration variables any pdf can be extracted
   This yields an indirect determination of any interesting quantity



# The Scanning Method

- The basis is the original approach by M.H. Schune & S. Plaszczynski used for the BABAR physics book
- The fit method was extended to include over 250 single measurements
- The four QCD parameters  $B_k$ ,  $f_B$ ,  $B_B$ ,  $\xi$  and  $V_{ub}$ ,  $V_{cb}$ , have significant theory uncertainties, thus they are scanned in the following way
  - We express each parameter in terms of  $x_0 \pm \sigma_x \pm \delta_x$ , where  $\sigma$  is a statistical uncertainty and  $\delta_x$  is the theory uncertainty
  - We select a specific value  $x^* \in [x_0 \delta_x, x_0 + \delta_x]$  as a model
  - We consider all models inside the  $[x_0 \delta_x, x_0 + \delta_x]$  interval
  - In each model the uncertainty  $\sigma_q$  is treated in a statistical way
- The uncertainties in the QCD parameters  $\eta_{cc}$ ,  $\eta_{ct}$ ,  $\eta_{tt}$ , and  $\eta_B$  and the quark masses  $m_c(\overline{m}_c)$  and  $m_b(\overline{m}_b)$  are treated like statistical uncertainties, since these uncertainties are relatively small → however, if necessary, we can scan over any of these parameters



# The Scanning Method

- We perform maximum likelihood fits using a frequentist approach
- A model is considered consistent with data if  $P(\chi^2_M)_{min} > 5\%$
- For consistent models we determine the best estimate and plot a 95% CL (ρ, η) contour → we overlay contours of consistent models
   → however, though only one of the contours is the correct one, we do not know which and thus show a representative numebr of them
- For accepted fits we also study the correlations among the theoretical parameters extending their range far beyond the range specified by the theorists
- We can input  $\alpha$ ,  $\phi_2$  and  $\gamma$ ,  $\phi_3$  via a likelihood function or directly using individual  $B \rightarrow \pi\pi$ ,  $\rho\pi$ ,  $\rho\rho$ ,  $a_1\pi$ ,  $b_1\pi$  measurements and GLW, ADS and Dalitz plot measurements in  $B \rightarrow D^{(*)}K^{(*)}$  & sin(2 $\beta$ + $\gamma$ ), respectively
- $\clubsuit$  We can further determine PP, PV VV amplitudes and strong phases using Gronau and Rosner parameterizations in powers of  $\lambda$

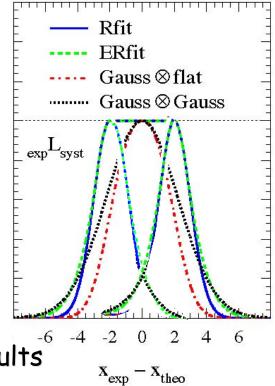


Work is in progress to include cos  $2\beta$ ,  $\beta_s$ ,  $A^q_{SL}$ ,  $\Delta\Gamma_s$  and  $\tau_s$ add contours of sin  $2\alpha$ ,  $\gamma$  and sin  $(2\beta+\gamma)$  and improve on display G. Eigen, Vxb workshop, SLAC, October 31 2009

# Differences among the 3 Methods

Likelihood

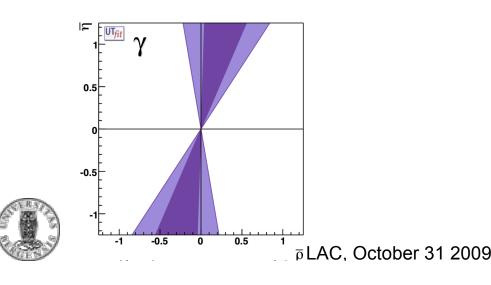
- Fit methodologies differ: 2 frequentist approaches vs 1 Bayesian approach
- Theory uncertainties are treated differently in the global fits
- Presently, measurement input values differ plus some assumptions differ
- For V<sub>ub</sub> and V<sub>cb</sub> there is an issue how to combine inclusive and exclusive results
  - Inclusive/exclusive averages  $\leftrightarrow$  individual results
  - Resulting errors
- ➡ We need to standardize measurement inputs, QCD parameters (at least numerical values should agree) and assumptions

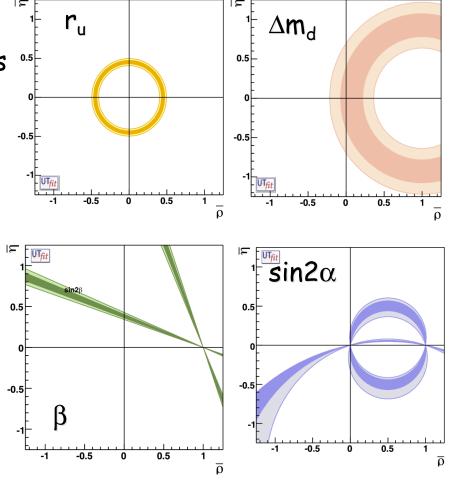




### Input Measurements from B Factories

- V<sub>ub</sub> and V<sub>cb</sub> measured in exclusive and inclusive semileptonic B decays
- $\Delta m_d$  from  $B_d B_d$  oscillations
- CP asymmetries  $a_{cp}(\psi K_S)$ from  $B \rightarrow ccK_s$  decays  $\rightarrow$  angle  $\beta$
- $\alpha$  from  $B \rightarrow \pi\pi$ ,  $B \rightarrow \rho\rho$ , &  $B \rightarrow \rho\pi CP$ measurements, add  $B \rightarrow a_1\pi$ ,  $B \rightarrow b_1\pi$

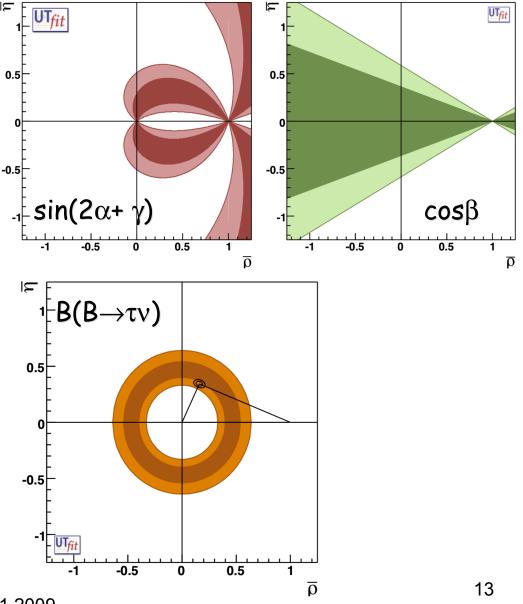




 $\blacksquare$  GLW, ADS and GGSZ analyses in  $B{\rightarrow}D^{(*)}K^{(*)}$ 

# Input Measurements from B Factories

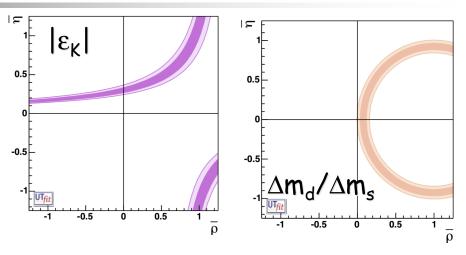
- sin(2 $\beta$ + $\gamma$ ) measurement from  $B \rightarrow D^{(*)}\pi(\rho)$
- cos 2  $\beta$  from B $\rightarrow$ J/ $\psi$ K<sup>\*</sup> and B $\rightarrow$ D<sup>0</sup> $\pi^{0}$
- $B \rightarrow \tau v$  branching fraction

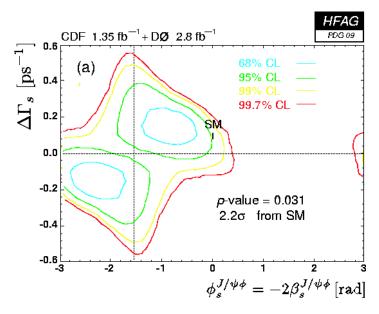




## Other Input Measurements

- $|\epsilon_{\rm K}|$  from CP violation in K decays
- $\Delta m_d / \Delta m_s$  from  $B_d B_d$  and  $B_s B_s$  oscillations
- CKM elements V<sub>ud</sub>, V<sub>us</sub>, V<sub>cd</sub>, V<sub>cs</sub>, V<sub>tb</sub>





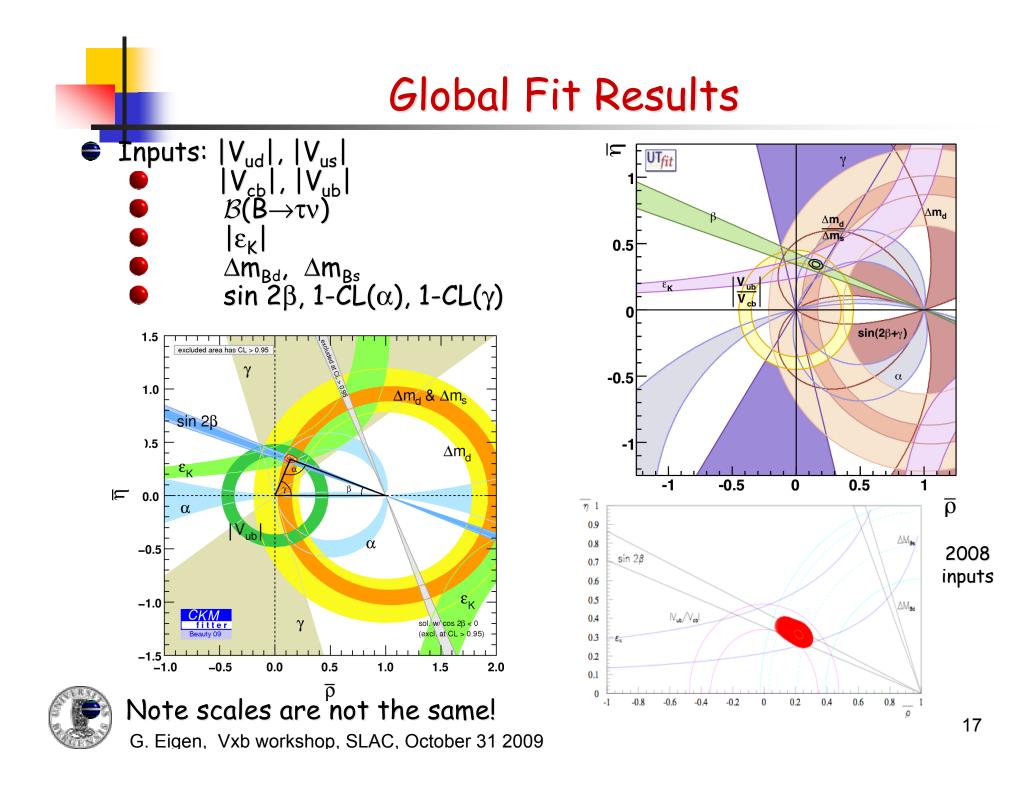


#### Measurement Inputs

	Observable	CKMfitter	UTfit		Scanning M	
•	V <sub>us</sub>	0.2246±0.0012	0.2259±0.0009		0.2258±0.0021	
	V <sub>ub</sub>   [10 <sup>-3</sup> ]	3.79±0.09±0.41*	<b>4.11</b> <sup>+0.27</sup> -0.28	(inc)		
			3.38±0.36	(exc)	3.84±0.16±0.29	(ex
	V <sub>cb</sub>   [10 <sup>-3</sup> ]	40.59±0.37±0.58*	41.54±0.73	(inc)		
			38.6±1.1	(exc)	40.9±1.0±1.6 (ex	(c)
	B(B→τν) [10 <sup>-4</sup> ]	1.73±0.35	1.51±0.33		1.79±0.72	
	∆m <sub>Bd</sub> [p <i>s</i> <sup>-1</sup> ]	0.507±0.005	0.507±0.005 17.77±0.12 2.229±0.010 0.671±0.023		0.508±0.005	
	∆m <sub>Bs</sub> [ps <sup>-1</sup> ]	17.77±0.12			17.77±0.12	
	ε <sub>κ</sub>   [10 <sup>-3</sup> ]	2.229±0.010			2.232±0.007	
	sin 2β	0.671±0.023			0.68±0.025	
	<b>α[ ππ, ρπ, ρρ]</b>	1-CL(α)	$-2\Delta ln(\mathcal{L})$		<b>Β, Տ, C for</b> ππ & ρρ	
	$\gamma$ [GGSZ, GLW, ADS]	1-CL(γ)	-2∆ln( <i>L</i> )		GGSZ, GLW, ADS	
	cos 2β	J/ψK*	J/ψK*, D <sup>ο</sup> π <sup>ο</sup>		To be done	
	2β+γ	D <sup>(*)</sup> π(ρ)	-2∆ln( <i>L</i> )		D <sup>(*)</sup> π(ρ)	
<u>I</u>	J	* use average values 15				

# Lattice QCD Inputs

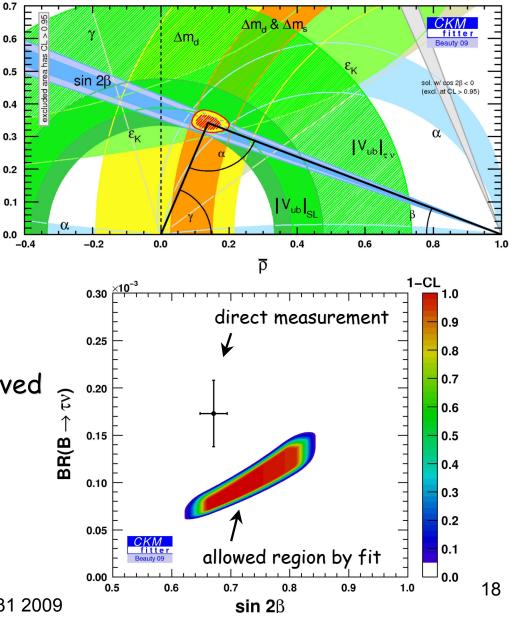
🖨 Parameter	CKMfitter			UTfit		Scanning	
	Mean	$\sigma_{\text{stat}}$	$\delta_{ ext{theo}}$	Mean	σ	Mean $\sigma_{\text{stat}}$ $\delta_{\text{theo}}$	
f <sub>Bs</sub> [f <sub>Bd</sub> ]	228 ±	<u>-</u> 3	±17	245	±25	[216 ±10 ± 20]	
f <sub>Bs</sub> ∕f <sub>Bd</sub>	1.199 ±	800.0	±0.023	1.21	±0.04		
B <sub>Bs</sub> [B <sub>Bd</sub> ]	1.23 ±	0.03	±0.05	1.22	±0.12	[1.29 ±.05 ±.08]	
Β <sub>Β</sub> ₅/Β <sub>Β</sub> ₄ [ξ]	1.05 ±	0.02	±0.05	1.00	±0.03	[1.2 ±.028 ±.05]	
B <sub>K</sub> [2 GeV]	0.525 ±	0.0036	±0.052				
Β <sub>κ</sub>	0.721	±0.005	±0.040	0.75	±0.07	0.79 ±0.04 ±0.09	
m <sub>c</sub> (m <sub>c</sub> ) [GeV]	1.286	±0.013	±0.040	1.3	±0.1	1.27±0.11	
$m_t(\overline{m}_t) [GeV]$	165.02 ±	1.16	±0.11	161.2	±1.7	163.3 ±2.1	
$\eta_{cc}$	$\eta_{cc}$ Calculated from $m_c(m_c) \& \alpha_s$		1.38	±0.53	1.46±0.22		
$\eta_{ct}$	η <sub>ct</sub> 0.47±0.04			0.47	±0.04	0.47±0.04	
η <sub>tt</sub> 0.5765±0.0065				0.574	±0.004	0.5765±0.0065	
η <sub>β</sub> (MS)	0.551±0.007			0.55	±0.01	0.551±0.007	
C S	0.1176±0.	.0020		0.119	±0.03	0.118	



# Tension from $\mathcal{B}(B \rightarrow \tau v)$

Ц

- $\mathcal{B}(B \rightarrow \tau v)$  is proportional to  $|V_{ub}|^2$  and  $f_{Bd}^2$
- from global fit
  B(B→τν)=(0.79<sup>+0.016</sup><sub>-0.010</sub>)×10<sup>-4</sup>
- WA:  $\mathcal{B}(B \to \tau v) = (1.73 \pm 0.35) \times 10^{-4}$
- 2.4σ discrepancy
- If B(B→τν) or sin 2β are removed χ<sup>2</sup><sub>min</sub> in global fit drops by 2.4 |V<sub>ub</sub>|, |V<sub>cb</sub>| remain unaffected





#### Constraints in the $m_H$ -tan $\beta$ Plane From BABAR/Belle average m<sub>B</sub><sup>2</sup> B<sup>+</sup> $\frac{m_{B}^{2}}{m_{H^{+}}^{2}(1+\varepsilon_{0}\cdot\tan\beta)}\tan^{2}\beta$ r<sub>\_H</sub> = | 1 we extract H⁺ B<sup>+</sup> $r_{_{_{H}}} = 1.67 \pm 0.34_{_{exp}} \pm 0.36_{_{f_{_{B}}},V_{_{ub}}}$ 95% C.L. exclusions 1000 10<sup>9</sup> ∆a<sub>µ</sub><4.6 We can use the 95% CL MH to present exclusions at 800 Dark matter 95% CL in the $m_{\mu_{\perp}}$ -tan $\beta$ plane r<sub>H</sub> 600 5 4 400 3 95%CL $B \rightarrow X_s \gamma$ 2 Κ→μν 1 200 atror 0 0.0 0.1 0.20.3 0.4 tan $\beta/m_{H}$ 0 ATLAS 50 discovery curve 10 0 20 30 40 50 60

9

tan β



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#### Model-Independent Analysis of UT

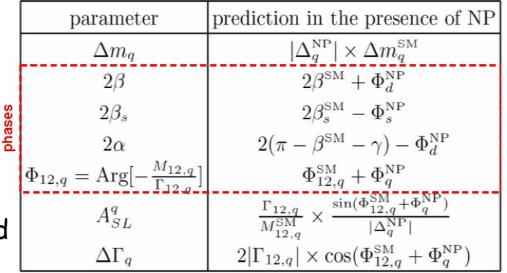
Solution Assume that new physics only affects short-distance part of  $\Delta B$ =2

• We use model-independent parameterization for  $B_d$  and  $B_s$ 

$$\frac{\left\langle \boldsymbol{B}_{q}^{0} \middle| \boldsymbol{\mathcal{H}}_{\Delta B=2}^{\textit{full}} \middle| \overline{\boldsymbol{B}}_{q}^{0} \right\rangle}{\left\langle \boldsymbol{B}_{q}^{0} \middle| \boldsymbol{\mathcal{H}}_{\Delta B=2}^{\textit{SM}} \middle| \overline{\boldsymbol{B}}_{q}^{0} \right\rangle} = \Delta_{q}^{\textit{NP}} = \left| \Delta_{q}^{\textit{NP}} \middle| \exp\left\{ 2i\phi_{q}^{\textit{NP}} \right\}$$

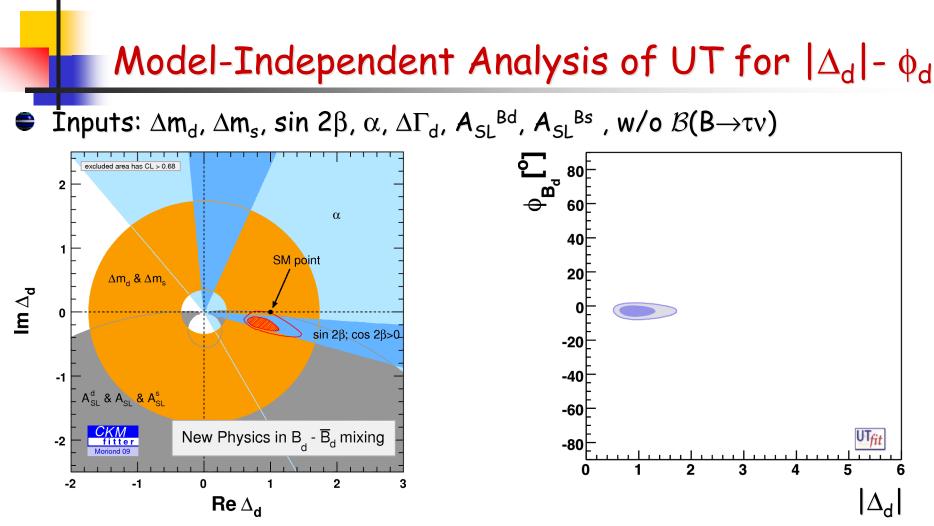
where  $H^{\text{full}} = H^{\text{SM}} + H^{\text{NP}}$ 

- Several observables are modified by the magnitude or phase of  $\Delta_q^{NP}$
- In  $B_d$  system we compare  $R_u \& \gamma$  with sin  $2\beta$ , sin $2\alpha$  and  $\Delta m_d$  that may be modified by NP parameters ( $\Delta_d$ ,  $\phi_d$ )





In B<sub>s</sub> system we compare R<sub>u</sub> &  $\gamma$  with  $\Delta m_s$ ,  $\beta_s$ , and  $\Delta \Gamma_s$  that may be modified by NP parameters ( $\Delta_d$ ,  $\phi_d$ )



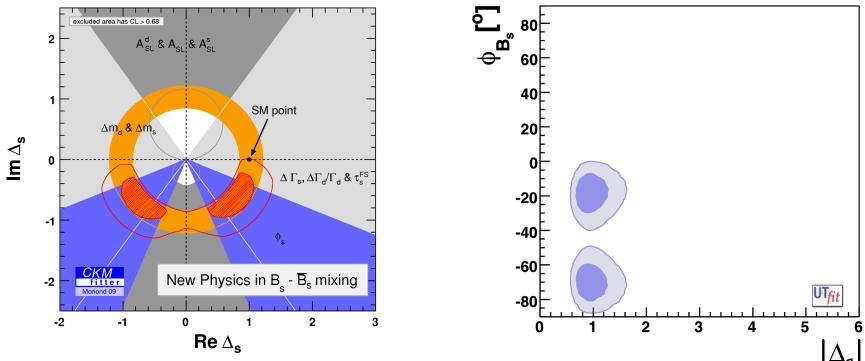
- Dominant constraints come from  $\beta$  and  $\Delta m_d$
- $\clubsuit$  Semileptonic asymmetries  $A_{SL}$  exclude symmetric solution with  $\eta \mbox{-}0$



 $\Delta_{d}=1 \text{ (SM) is disfavored by 2.1}\sigma \text{ (discrepancy }\mathcal{B}(B\to\tau\nu) \text{ and sin 2}\beta)$   $\Rightarrow \phi_{d}^{NP}=(-12^{+9}-_{6})^{0} @ 95\% CL \text{ (}\Rightarrow \text{discrepancy is 0.6}\sigma \text{ w/o }\mathcal{B}(B\to\tau\nu)\text{ )}$ G. Eigen, Vxb workshop, SLAC, October 31 2009 21

#### Model-Independent Analysis of UT for $|\Delta_s|$ - $\phi_s$

 $= \text{Inputs: } \phi_s, \Delta m_d, \Delta m_s, A_{SL}^{Bd}, A_{SL}^{Bs}, \Delta \Gamma_s, \tau_s, \mathcal{B}(B \rightarrow \tau v)$ 



- Dominant constraints come from direct measurements of  $\phi_s$ ,  $\Delta \Gamma_s^{1/2}$ in B $\rightarrow$ J/ $\psi \phi$  and  $\Delta m_s$  from the Tevatron
- $\phi_s$  is 2.2 $\sigma$  away from the SM prediction



 $\Delta_s$  =1 is disfavored at 1.9 $\sigma$  level independent of  $\mathcal{B}(B \rightarrow \tau v)$ 

# Final Remarks

- Among the 3 global CKM fitting methods we need to standardize
  - All measurement inputs
  - What QCD parameters to use, their central values, their statistical errors and their theory errors
  - The notation for quantities in the text, on plots and in equations
- We need to specify which input parameters are used in the fits and standardize on the assumptions
  - $\rightarrow$  This is important for comparing results
- We will present results in form of plots with values listed in tables
   we accompany the results with a few remarks
   in particular in cases of discrepancies we need to discuss them
- Most of the writing probably has to be done by the co-conveners



## More Recent Publications

- CKMfitter publications
  - J. Charles et al., Eur. Phys. J. C41, 1-135, 2005.

#### UTfit publications:

- M. Bona et al., JHEP 0610:081, 2006.
- M. Bona et al., Phys.Rev.D76:014015, 2007.
- M. Bona et al., Phys.Rev.Lett.97:151803, 2006.
- M Bona et al., Phys.Lett.B687:61-69, 2010.

#### Scanning method

- G. Eigen et al., Eur.Phys.J.C33:S644-S646,2004.
- G.P. Dubois-Felsmann et al., hep-ph/0308262.

