

Rencontres de physique
de l'infiniment grand à l'infiniment petit

27 juillet 2011

Cent ans après Rutherford

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Institut de Physique Nucléaire
CNRS/IN2P3, Université Paris-Sud 11

En 1911

Rutherford introduit un modèle d'atome

pour expliquer

les résultats expérimentaux de Geiger et Marsden

de 1909

Atomisme au dix neuvième siècle

Chimie

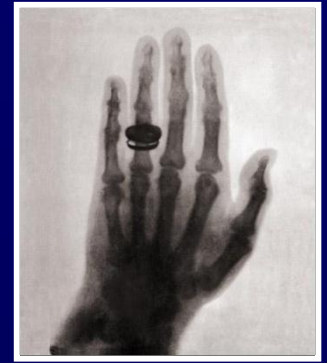
Théorie cinétique des gaz

1895 - Rayons de Röntgen (Rayons X)

1896 - Rayons uraniques de Becquerel (Radioactivité)

1897 - Corpuscules d'électricité par J.J. Thomson (Electrons)

1895 - Rayons de Röntgen (Rayons X)



Contexte

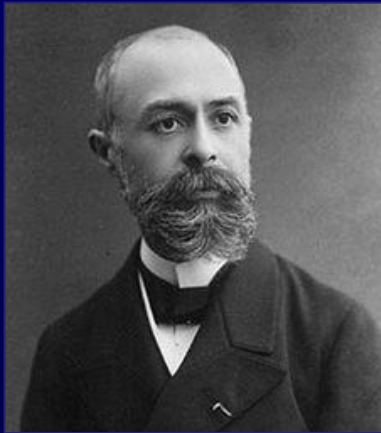
Etude des
décharges électriques
dans les gaz



Tube de Crookes

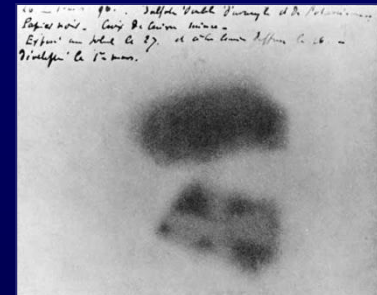
1897 - Corpuscules d'électricité par J.J. Thomson (Electrons)

1896 - Rayons uraniques de Becquerel (Radioactivité)



Contexte

Etude de la
phosphorescence
des sels d'uranium



Travaux développés par
Pierre et Marie Curie

Découverte
du Polonium et du Radium



L'atome en 1900



L'atome en 1900



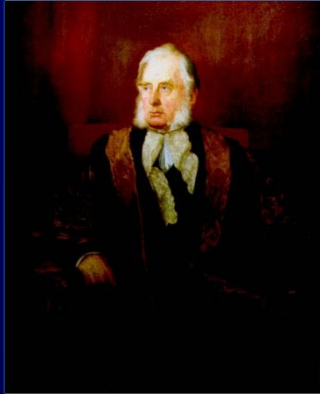
Congrès International de Physique

- | | |
|--|--------------------|
| 1. Questions générales. Unités. Mesures | 7, 9 et 11 août |
| 2. Physique mécanique et moléculaire | 7, 9 et 11 août |
| 3. Optique et thermodynamique | 8 et 10 août |
| 4. Electricité et magnétisme | 8, 10 et 11 août |
| 5. Magnéto-optique. Rayons cathodiques. Rayonnement de l'uranium. Etc. | 7, 8, 9 et 10 août |
| 6. Physique cosmique | 7, 8, 9 et 10 août |
| 7. Physique biologique | 7, 8 et 9 août |

Séance générale de clôture
12 août



Cavendish Laboratory (Cambridge)



Financé par
W. Cavendish
(7^{ème} duc of Devonshire)

1874

Directeur
Professeur de
physique expérimentale



1874 - 1879

1879 - 1884

1884 - 1919

1919 - 1937

1938 - 1954

J. C. Maxwell

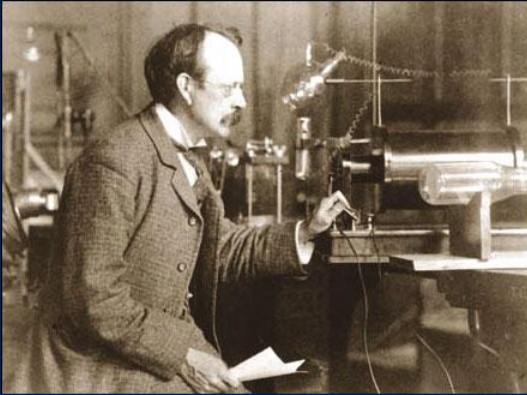
Lord Rayleigh

J.J. Thomson

E. Rutherford

W.L. Bragg

Joseph John Thomson



- 1856 Naissance à Manchester
- 1870 Etudes au Owens Collège (Manchester)
- 1876 Boursier au Trinity Collège (Cambridge)
- 1880 "Wrangler". Reçu 2^{ème} au "Tripos" de mathématiques de Cambridge

1884 - 1919
Directeur du
Laboratoire Cavendish

1906

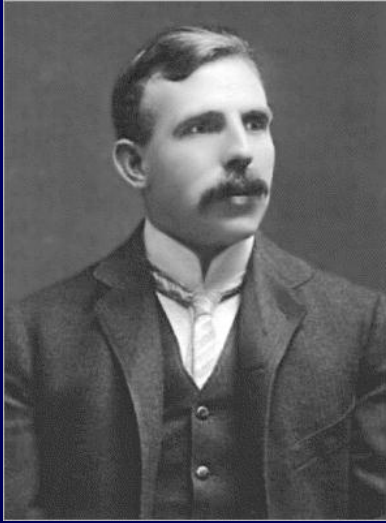
Prix Nobel de physique

En témoignage des grands mérites de ses recherches théoriques et expérimentales sur la conduction de l'électricité dans les gaz



1940 Décès à Cambridge

Ernest Rutherford



1871 Naissance à Nelson (Nouvelle Zélande)
Boursier au Nelson puis au Canterbury College
de Christchurch (N.Z.). B.A., M.A et B.Sc.

1895 Bourse pour le Cavendish de Cambridge

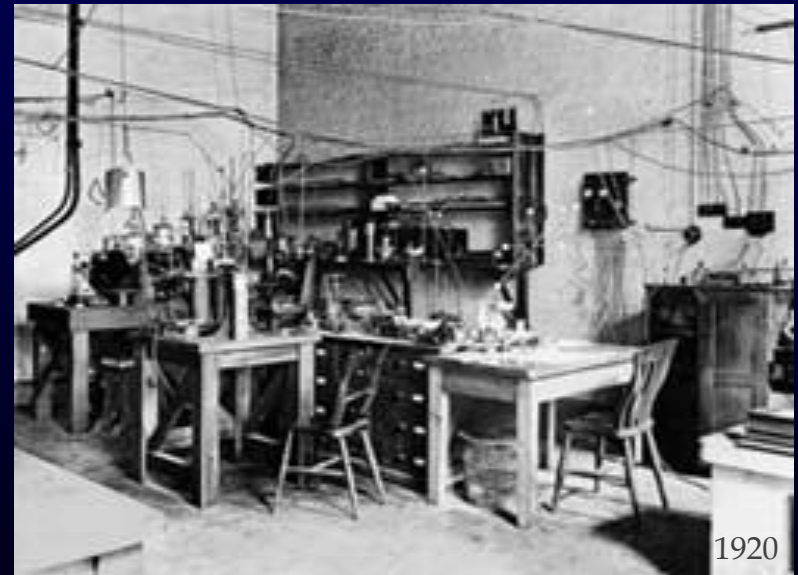
1898 Professeur à McGill, Montreal (Canada)

1908 Prix Nobel de chimie pour ses recherches sur la
désintégration des éléments et la chimie des substances radioactives

1907 Université de Manchester



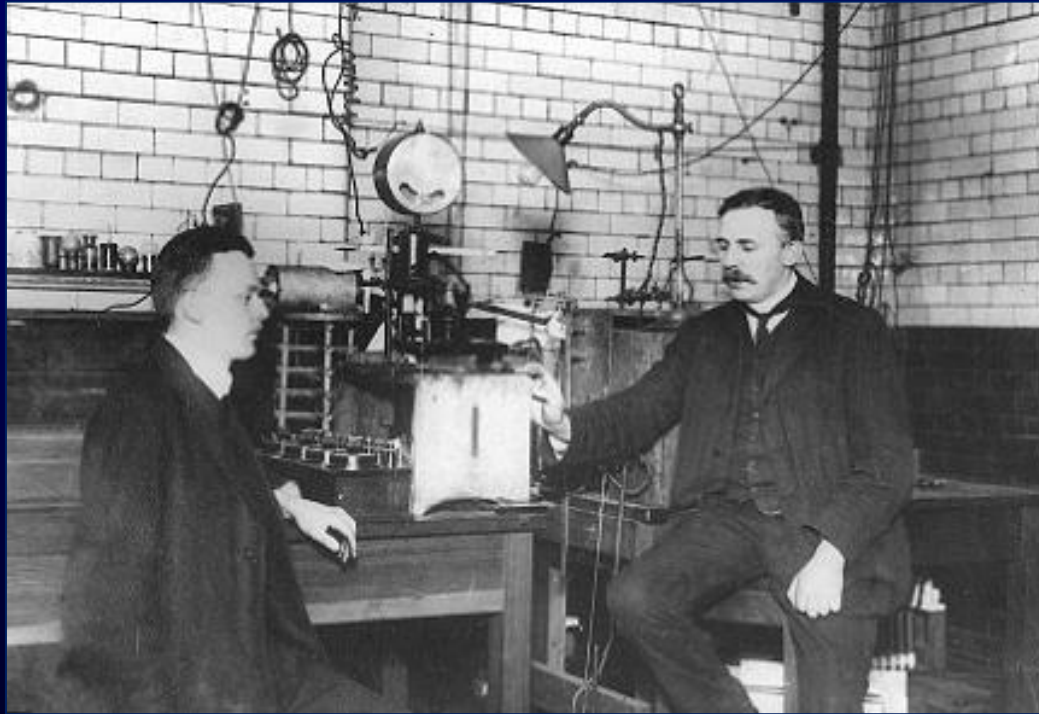
1919 Cavendish Laboratory



1937 Décès à Cambridge

Les expériences de Geiger et Marsden

1909 On a Diffuse Reflexion of the α -Particles



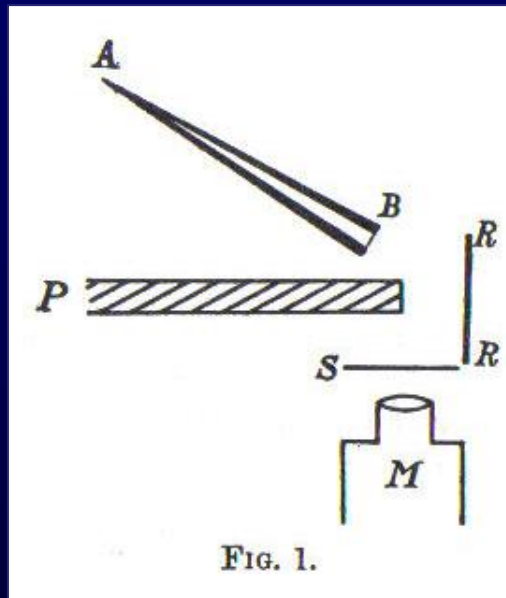
Les expériences de Geiger et Marsden

1909 On a Diffuse Reflexion of the α -Particles

Source gazeuse
(Tube conique AB)

Plaque de plomb B

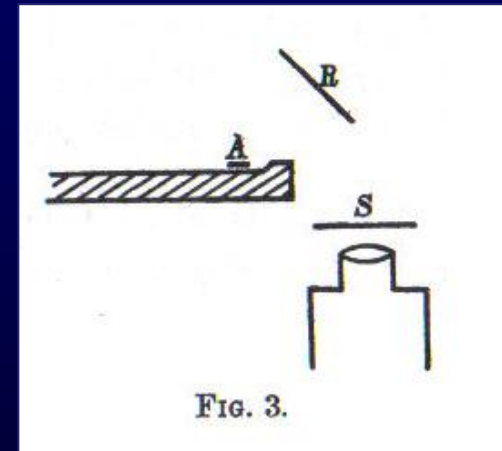
Détecteur S
Microscope M



Source
en A

Réflecteur
en platine R

Réflecteur
RR



1. Metal.	2. Atomic weight, A.	3. Number of scintillations per minute, Z.	4. A/Z.
Lead	207	62	30
Gold	197	67	34
Platinum.....	195	63	33
Tin	119	34	28
Silver	108	27	25
Copper.....	64	14.5	23
Iron	56	10.2	18.5
Aluminium.....	27	3.4	12.5

Le modèle de Rutherford

1911 The Scattering of the α and β Particles by Matter and the Structure of the Atom


Theory of structure of atom

Suppose atom consists of + charge Ze at centre + - charge electrons distributed throughout sphere of radius a .

Force at P on electron = $Ne^2 \left\{ \frac{1}{r^2} - \frac{4}{3a^3} \frac{1}{r^3} \right\}$

$$= Ne^2 \left\{ \frac{1}{r^2} - \frac{4}{3a} \right\} = \neq \neq$$

Suppose charged particles e mass m moves through atom so that deflection is small but L^2 distance from centre = a



deflection force L^2 double bracket at P

$$= Ne^2 \left\{ \frac{1}{r^2} - \frac{4}{3a} \right\} \cos \theta$$

\therefore total L^2 double bracket = $dd = \frac{Ne^2}{m} \left\{ \frac{1}{r^2} - \frac{4}{3a} \right\} \frac{a}{r}$

\therefore Hence u acquired in passing through atom L^2 double bracket

$$u = \int dd \cdot dt = Ne^2 ka \cdot \frac{ds}{v^2}$$

$$= \frac{Ne^2}{mv} \int \left(\frac{1}{r^2} - \frac{4}{3a} \right) \frac{a}{r} \cdot \frac{r dr}{r^2 - a^2}$$

$$= \frac{2Ne^2}{mv} \int \frac{a(1/r^3 - 4/3a)}{r^2 - a^2} \cdot \frac{dr}{2}$$

$$= \frac{Ne^2}{mv} \left(\frac{1}{2a} \ln \frac{r+a}{r-a} - \frac{4}{3a} \ln \frac{r+a}{r-a} \right) \cdot \frac{dr}{2}$$

[669]

LXXIX. *The Scattering of α and β Particles by Matter and the Structure of the Atom.* By Professor E. RUTHERFORD, F.R.S., University of Manchester*.

§ 1. IT is well known that the α and β particles suffer deflexions from their rectilinear paths by encounters with atoms of matter. This scattering is far more marked for the β than for the α particle on account of the much smaller momentum and energy of the former particle. There seems to be no doubt that such swiftly moving particles pass through the atoms in their path, and that the deflexions observed are due to the strong electric field traversed within the atomic system. It has generally been supposed that the scattering of a pencil of α or β rays in passing through a thin plate of matter is the result of a multitude of small scatterings by the atoms of matter traversed. The observations, however, of Geiger and Marsden † on the scattering of α rays indicate that some of the α particles must suffer a deflexion of more than a right angle at a single encounter. They found, for example, that a small fraction of the incident α particles, about 1 in 20,000, were turned through an average angle of 90° in passing through a layer of gold-foil about $\cdot 00004$ cm. thick, which was equivalent in stopping-power of the α particle to 1.6 millimetres of air. Geiger ‡ showed later that the most probable angle of deflexion for a pencil of α particles traversing a gold-foil of this thickness was about $6^\circ 87'$. A simple calculation based on the theory of probability shows that the chance of an α particle being deflected through 90° is vanishingly small. In addition, it will be seen later that the distribution of the α particles for various angles of large deflexion does not follow the probability law to be expected if such large deflexions are made up of a large number of small deviations. It seems reasonable to suppose that the deflexion through a large angle is due to a single atomic encounter, for the chance of a second encounter of a kind to produce a large deflexion must in most cases be exceedingly small. A simple calculation shows that the atom must be a seat of an intense electric field in order to produce such a large deflexion at a single encounter.

Recently Sir J. J. Thomson § has put forward a theory to

* Communicated by the Author. A brief account of this paper was communicated to the Manchester Literary and Philosophical Society in February, 1911.

† Proc. Roy. Soc. LXXXII. p. 495 (1909).

‡ Proc. Roy. Soc. LXXXIII. p. 482 (1910).

§ Camb. Lit. & Phil. Soc. xv. pt. 5 (1910).

Article lu le 7 mars 1911 à Manchester Literary and Philosophical Society

1911

Le modèle de Rutherford

[669]

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‡ Proc. Roy. Soc. lxxxiii. p. 492 (1910).

§ Camb. Lit. & Phil. Soc. xv. pt. 5 (1910).

1911

Le modèle de Rutherford

674 The probability m of entering an atom within a distance p of its centre is given by

$$m = \pi p^2 nt.$$

Chance dm of striking within radii p and $p + dp$ is given by

$$dm = 2\pi pnt \cdot dp = \frac{\pi}{4} nt^2 \cot \phi/2 \operatorname{cosec}^2 \phi/2 d\phi. \quad (2)$$

since $\cot \phi/2 = 2p/b$.

The value of dm gives the fraction of the total number of particles which are deviated between the angles ϕ and $\phi + d\phi$.

The fraction ρ of the total number of particles which are deflected through an angle greater than ϕ is given by

$$\rho = \frac{\pi}{4} nt^2 \cot^2 \phi/2. \quad (3)$$

The fraction ρ which is deflected between the angles ϕ_1 and ϕ_2 is given by

$$\rho = \frac{\pi}{4} nt^2 \left(\cot^2 \frac{\phi_1}{2} - \cot^2 \frac{\phi_2}{2} \right). \quad (4)$$

It is convenient to express the equation (2) in another form for comparison with experiment. In the case of the α rays, the number of scintillations appearing on a constant area of a zinc sulphide screen are counted for different angles with the direction of incidence of the particles. Let r = distance from point of incidence of α rays on scattering material, then if Q be the total number of particles falling on the scattering material, the number y of α particles falling on unit area which are deflected through an angle ϕ is given by

$$y = \frac{Qdm}{2\pi r^2 \sin \phi \cdot d\phi} = \frac{nt^2 \cdot Q \cdot \operatorname{cosec}^4 \phi/2}{16r^2}. \quad (5)$$

Since $b = \frac{2NeE}{mu^2}$, we see from this equation that the number of α particles (scintillations) per unit area of zinc sulphide screen at a given distance r from the point of

incidence of the rays is proportional to

- (1) $\operatorname{cosec}^4 \phi/2$ or $1/\phi^4$ if ϕ be small;
- (2) thickness of scattering material t provided this is small;
- (3) magnitude of central charge Ne ;
- (4) and is inversely proportional to $(mu^2)^2$, or to the fourth power of the velocity if m be constant.

In these calculations, it is assumed that the α particles scattered through a large angle suffer only one large deflexion. For this to hold, it is essential that the thickness of the scattering material should be so small that the chance of a second encounter involving another large deflexion is very small. If, for example, the probability of a single deflexion ϕ in passing through a thickness t is $1/1000$, the probability of two successive deflexions each of value ϕ is $1/10^6$, and is negligibly small.

The angular distribution of the α particles scattered from a thin metal sheet affords one of the simplest methods of testing the general correctness of this theory of single scattering. This has been done recently for α rays by Dr. Geiger*, who found that the distribution for particles deflected between 30° and 150° from a thin gold-foil was in substantial agreement with the theory. A more detailed account of these and other experiments to test the validity of the theory will be published later.

§ 4. Alteration of velocity in an atomic encounter.

It has so far been assumed that an α or β particle does not suffer an appreciable change of velocity as the result of a single atomic encounter resulting in a large deflexion of the particle. The effect of such an encounter in altering the velocity of the particle can be calculated on certain assumptions. It is supposed that only two systems are involved, viz., the swiftly moving particle and the atom which it traverses supposed initially at rest. It is supposed that the principle of conservation of momentum and of energy applies, and that there is no appreciable loss of energy or momentum by radiation.

* Manch. Lit. & Phil. Soc. 1910.

on e
with
 $p = 8$
initia

1911

Le modèle de Rutherford

1. Metal.	2. Atomic weight, A.	3. Number of scintillations per minute, Z.	4. A/Z.
Lead	207	62	30
Gold	197	67	34
Platinum.....	195	63	33
Tin	119	34	28
Silver	108	27	25
Copper.....	64	14·5	23
Iron	56	10·2	18·5
Aluminium	27	3·4	12·5

Résultats de Geiger et Marsden en 1909

Metal	Atomic weight	z	$z/A^{3/2}$
Lead	207	62	208
Gold	197	67	242
Platinum	195	63	232
Tin	119	34	226
Silver	108	27	241
Copper	64	14·5	225
Iron	56	10·2	250
Aluminium	27	3·4	243
Average			233

Repris par Rutherford en 1911

Conseil de Physique

Institut International de Physique Solvay



VERSCHAFFELT	LAUE	RUBENS	GOLDSCHMIDT	HERZEN	LINDEMANN	de BROGLIE	POPE	GRUNISEN	HOSTELET		
HASENOHRL	JEANS	BRADG	Mme CURIE	SOMMERFELD	EPSTEIN			KNUDSEN	LANGEVIN		
NERNST	RUTHERFORD	WIEN	J.J. THOMSON	WARBURG	LORENTZ	BILLOUIN	BARLOW	KAMERLINGH ONNES	WOOD	GOLY	WEISS

(27 -31 octobre 1913)

Niels Bohr



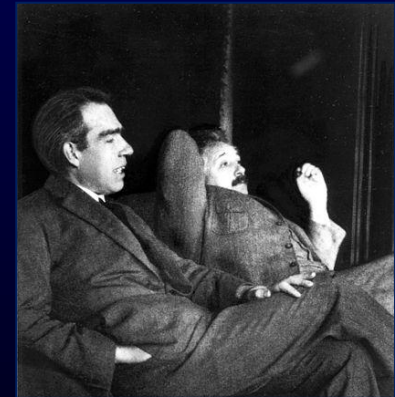
- 1885 Naissance à Copenhague
Fils de C. Bohr, professeur de médecine, recteur
- 1903 Etudes à l'Université de Copenhague
- 1911 Thèse (Copenhague) sur "la théorie électronique des métaux"



- 1916 Professeur à Copenhague
- 1921 Institut de physique théorique

1922
Prix Nobel
de physique

pour ses contributions à la
recherche de la structure
des atomes et sur le
rayonnement qu'ils émettent



1962 Décès à Copenhague

Niels Bohr



13 mai 1911 à Copenhague

Thèse : la théorie électronique des métaux



Début octobre 1911

Laboratoire Cavendish à Cambridge

16 mars au 3 mai 1912 à Manchester

Introduction aux méthodes
expérimentales en radioactivité



Fin juillet 1912 Mariage au Danemark



En juillet 1913, Professeur à Copenhague (Etudiants en médecine)

Sur proposition de Rutherford (mai 1914) ,
Professeur à Manchester (octobre 1914 à juillet 1916)

1913

L'atome de Bohr

On the Constitution of Atoms and Molecules

Publication en 3 parties dans
Philosophical Magazine

Partie 1

Envoyée par Rutherford le 5 avril 1913
Publiée en juillet 1913

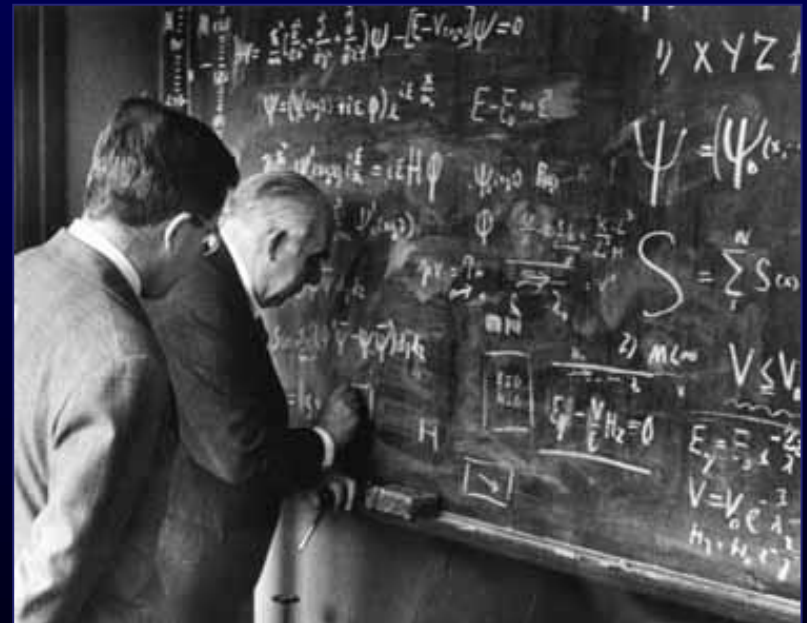
1914



1918

Thomson
Rutherford

Bohr



Bohr
1913



Rutherford
1911

L'activité rationaliste de la physique contemporaine
(Gaston Bachelard)