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A Short Introduction to spallation reactions.
Nuclear transport

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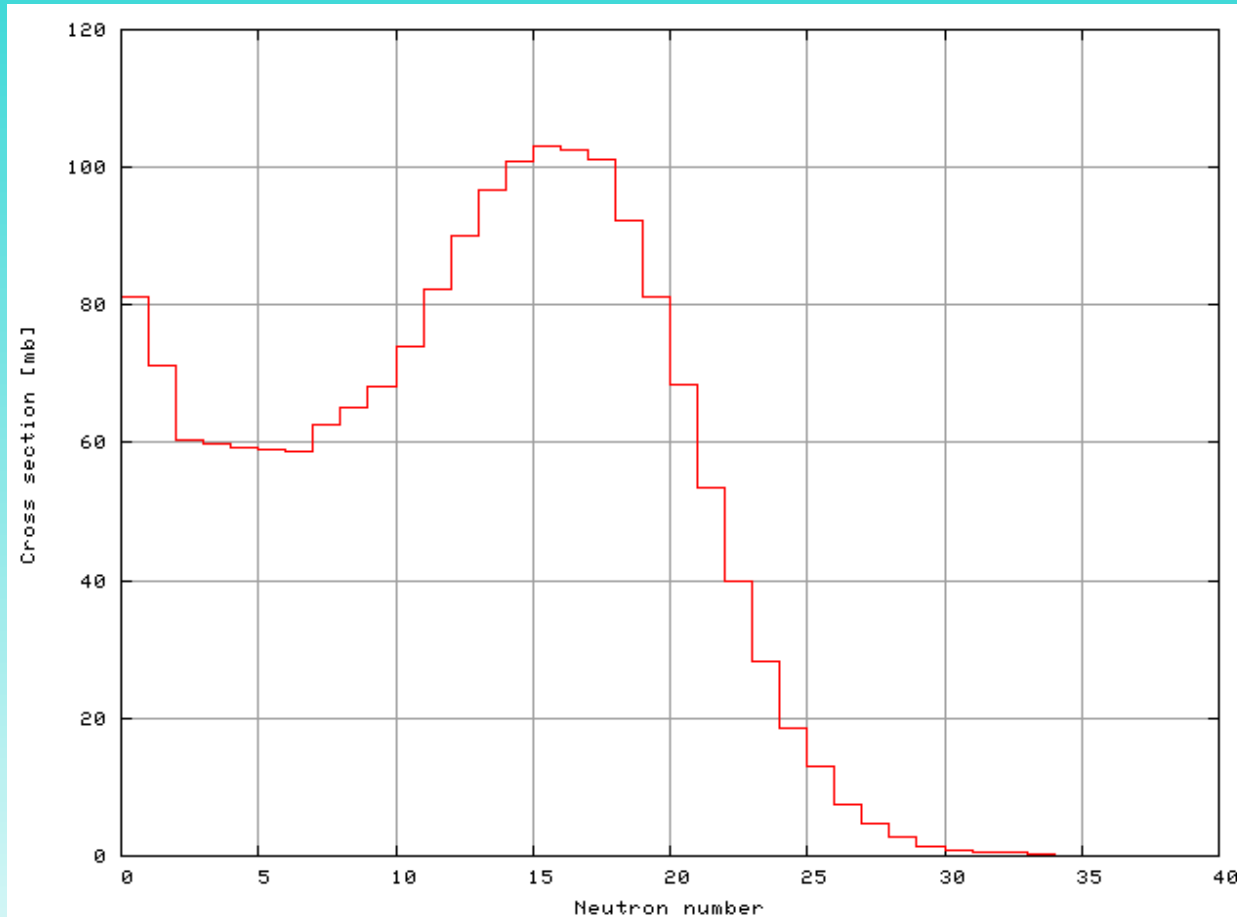


HLPW10, Spa, 6-8 March 2011

- *Spallation reactions*
- *Nuclear reactions as transport processes*
- *The collision regime. INC= a simple tool*
- *The validity of INC. Unification at intermediate energy*

SPALLATION REACTIONS

Many particles are expelled

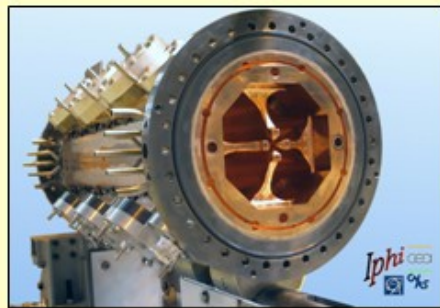
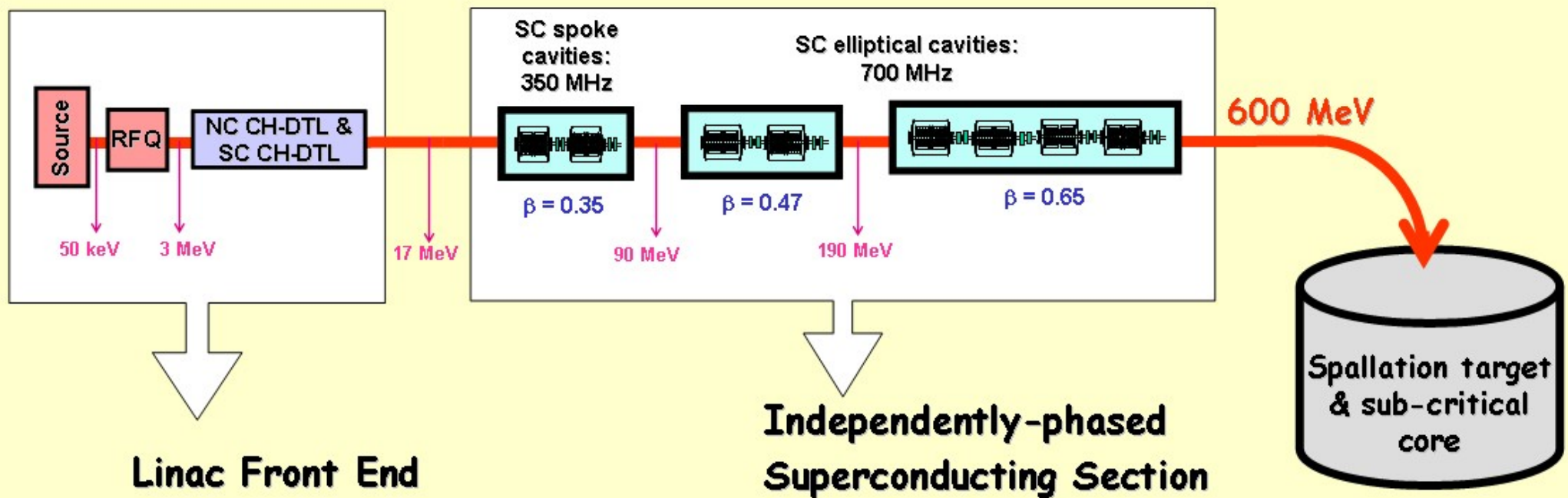


p (1.2 GeV) + Pb

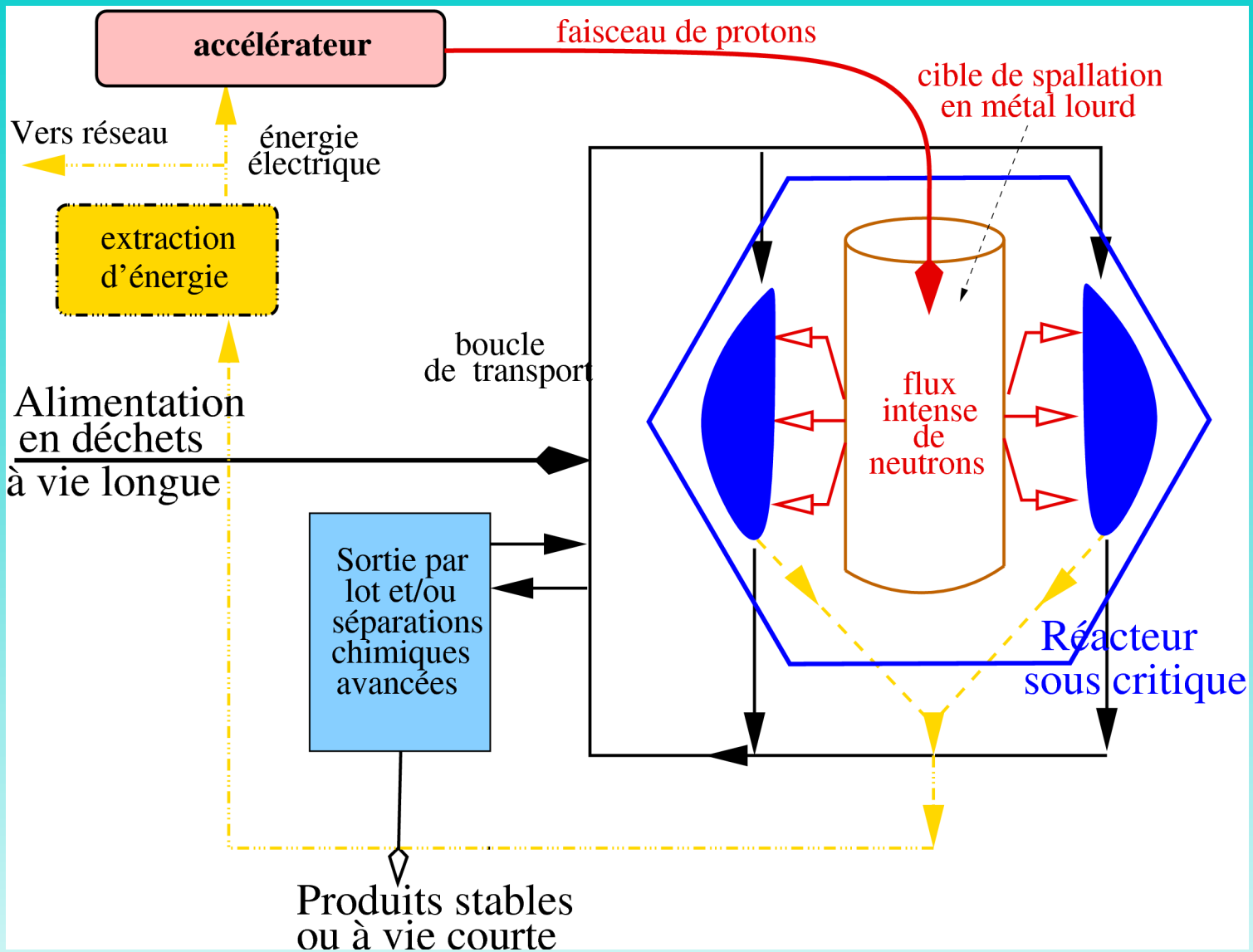
+ p, π , d, t, ... Spallation reaction=high-energy hadron-induced reactions

Macroscopic « spallation reactions » :

The MYRRHA accelerator reference scheme (2010)

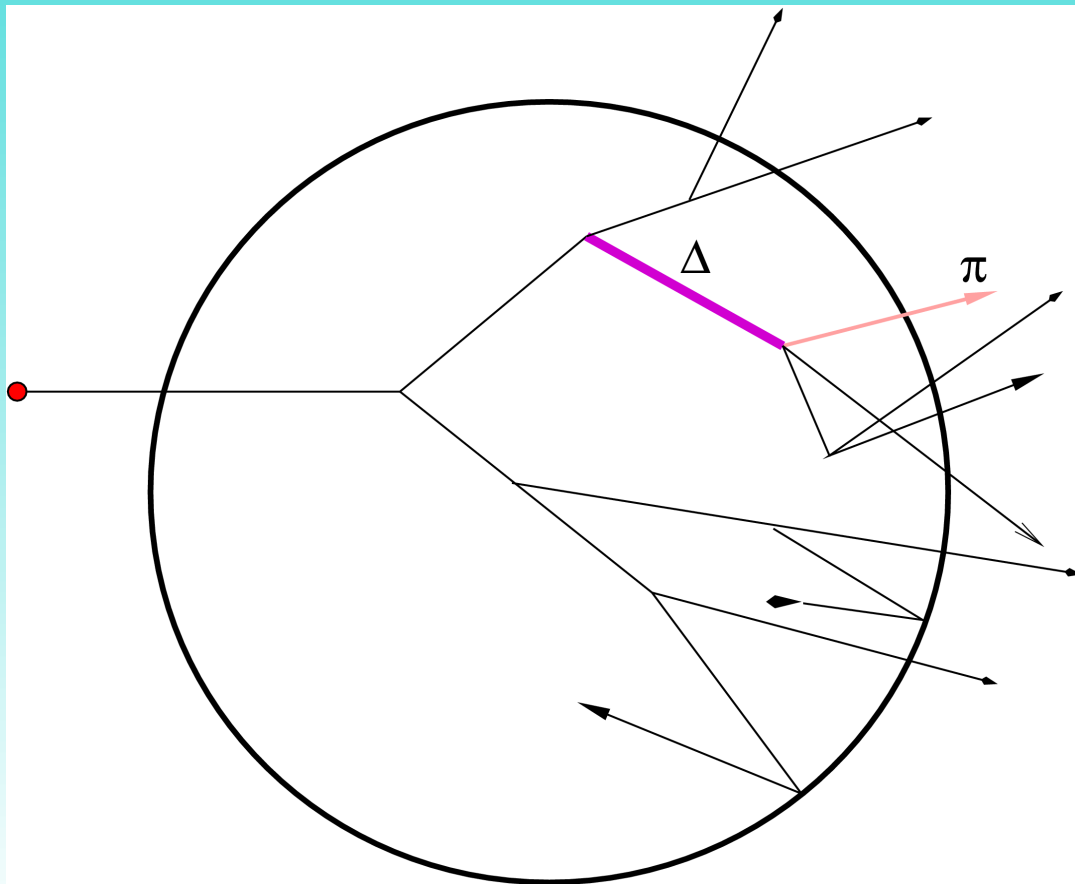


MYRRHA project is operated by SCK-CEN Mol



2. NUCLEAR REACTIONS AS TRANSPORT PROCESSES

At high energy, nuclear reactions are expected to proceed by binary collisions (R.Serber, 1947)



Reaction= diffusion of particles in the nuclear medium (under collisions, creation, destruction, modification ,...)

Heavy ion collisions= diffusion of projectile part. in the target + diffusion of target part. in the projectile

Diffusion processes:

- deterministic equations for average quantities
- simulations

The production of high-energy protons in central relativistic nuclear collisions*1

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Abstract

We extend our model of **transport theory** to be applicable to the inclusive production of protons with very high energy. We then consider the angular distribution of such protons, produced in a central collision of Ar on KC1 at 800 MeV per nucleon. The slight anisotropy observed in the data can be explained by a finite value of the friction constant which in turn determines the number of collisions needed for equilibrium to be reached. We also show that these data are **quite sensitive to the reaction geometry** and cannot be explained by the firestreak model.

Article Outline

- [References](#)

*1 Supported in part by the Bundesministerium für Forschung und Technologie.** Present address.

Nuclear Physics A Volume 360, Issue 2, 11 May 1981, Pages 435-443

3. THE COLLISION REGIME. INTRANUCLEAR CASCADE (INC) AS A TOOL

The nuclear transport equation :

Schrödinger many-body equation

$$\left(\frac{\partial}{\partial t} + \frac{\vec{p}}{m} \cdot \vec{\nabla} - \frac{1}{m} \vec{\nabla} U \cdot \vec{\nabla}_{\vec{p}}\right) f(\vec{r}, \vec{p}, t) =$$
$$\int \frac{d^3 p_2}{(2\pi)^3} \int \frac{d^3 p_3}{(2\pi)^3} \int \frac{d^3 p_4}{(2\pi)^3} |W(12 \rightarrow 34)|^2 \{f_3 f_4 (1-f_1)(1-f_2) - f_1 f_2 (1-f_3)(1-f_4)\} \delta(\vec{p}) \delta(e(p))$$

Approximations :

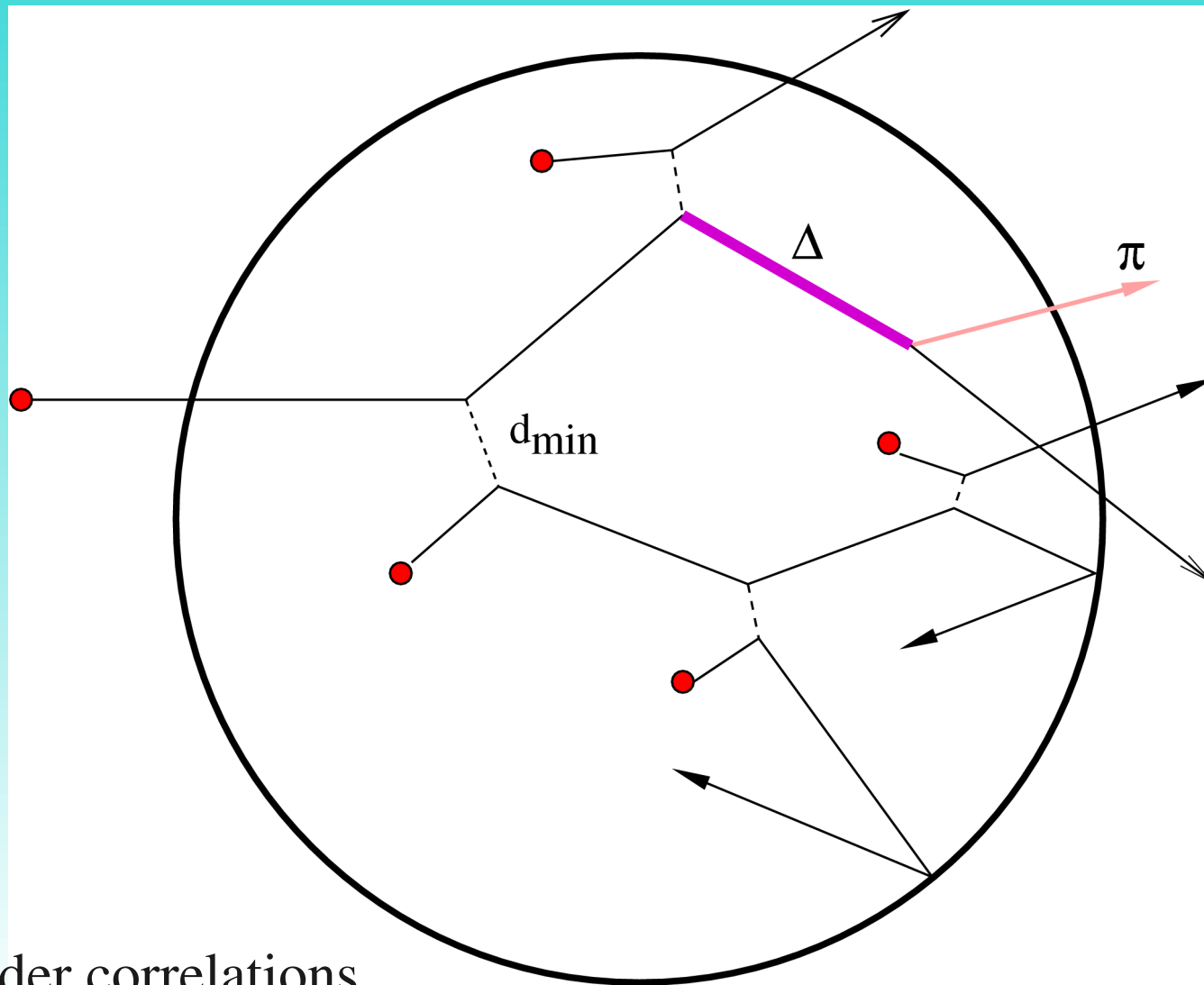
- closure (2 body correlations are unimportant)
- low-gradient expansion
- independence of collisions

$$f_i = f(\vec{r}_i, \vec{p}_i, t), f_1 = f$$

$$\lambda_B \ll t_{coll} v \ll d$$

For U fixed, and $W^2 = v_{12} \sigma$

Simulations are easy and correspond to INC



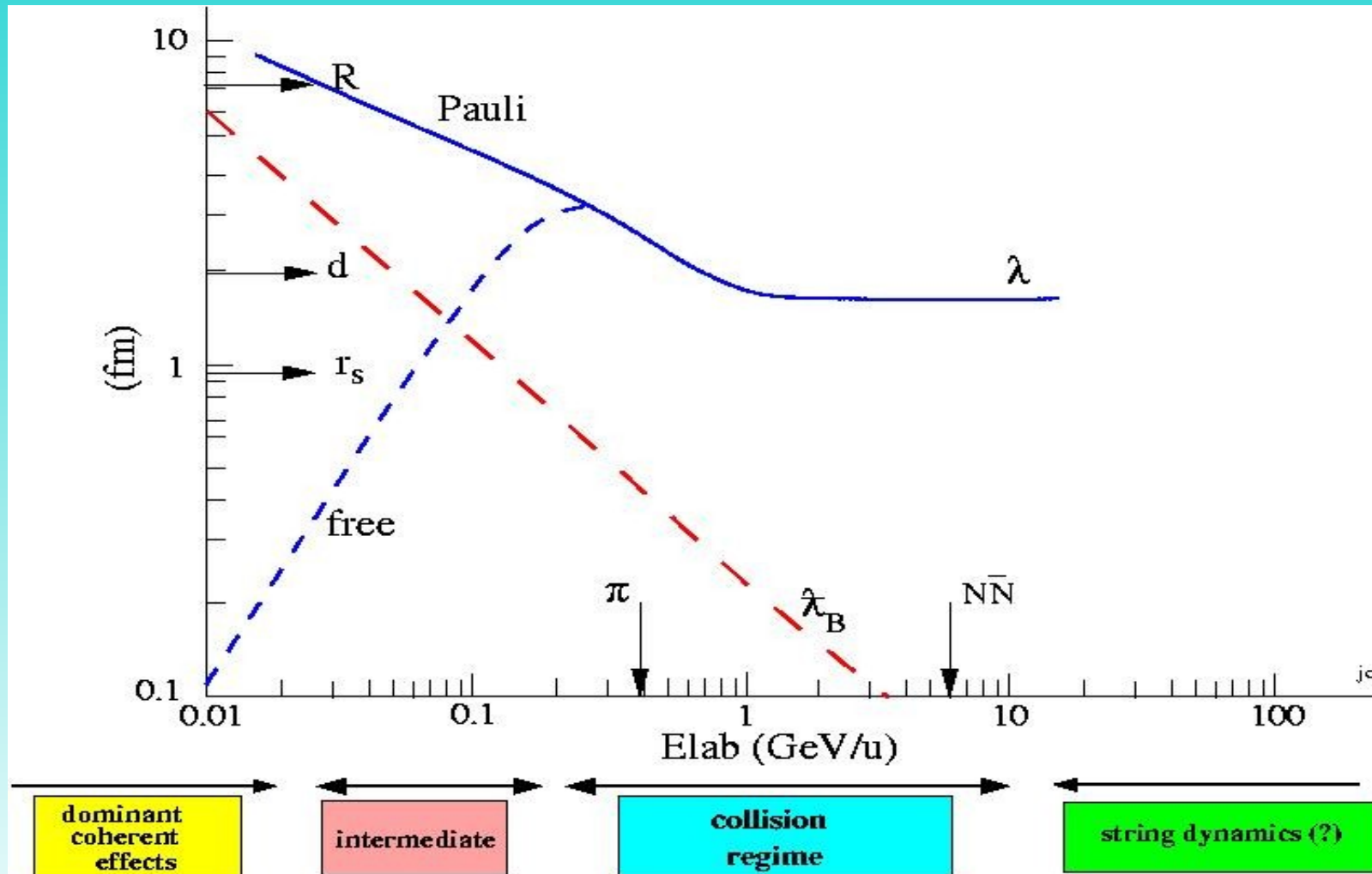
+higher-order correlations

4. VALIDITY OF INC. UNIFICATION AT INTERMEDIATE ENERGY

Low energy

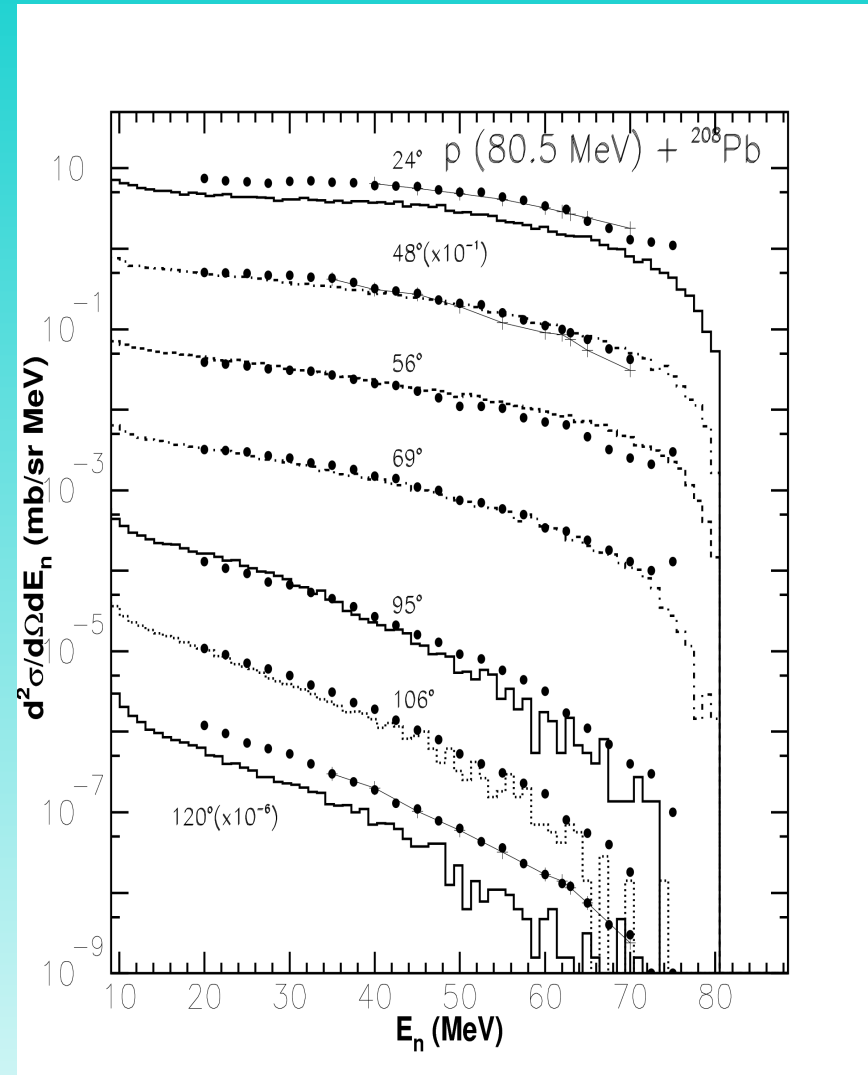
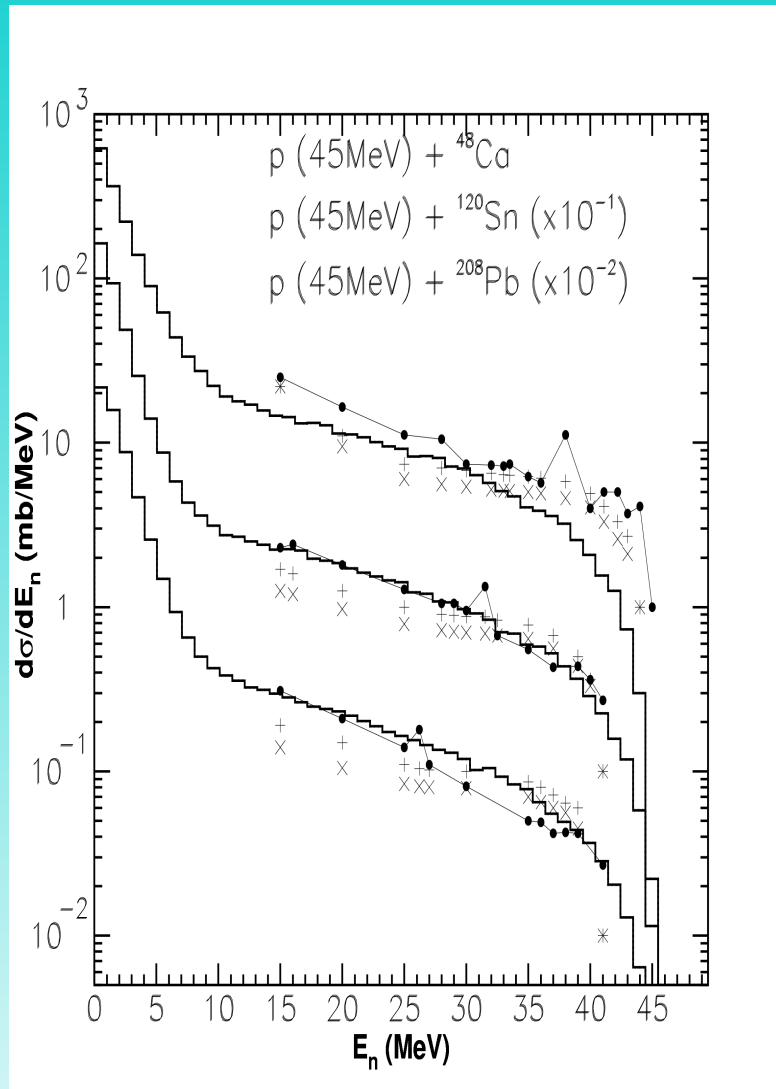
$$\lambda_B \ll t_{coll} v \ll d$$

→ $T_{lab} > 200 \text{ MeV}$



High energy → $T_{lab} < 20 \text{ GeV}$

For T_{lab} below 200 MeV, INC is working well



Explanation : Pauli blocking makes distance between collisions large enough

Typical alternative models= pre-equilibrium models

Collision process=particle-hole excitations, with various transition probabilities

Equivalent physics :

- probabilities are dictated by cross sections
- continuous phase space in INC = discrete p-h space with finite width

Interaction of an incident particle with a Fermi gas system dominated by « Pauli-blocked » binary collisions

CONCLUSIONS

Let us listen to the following lecturers

